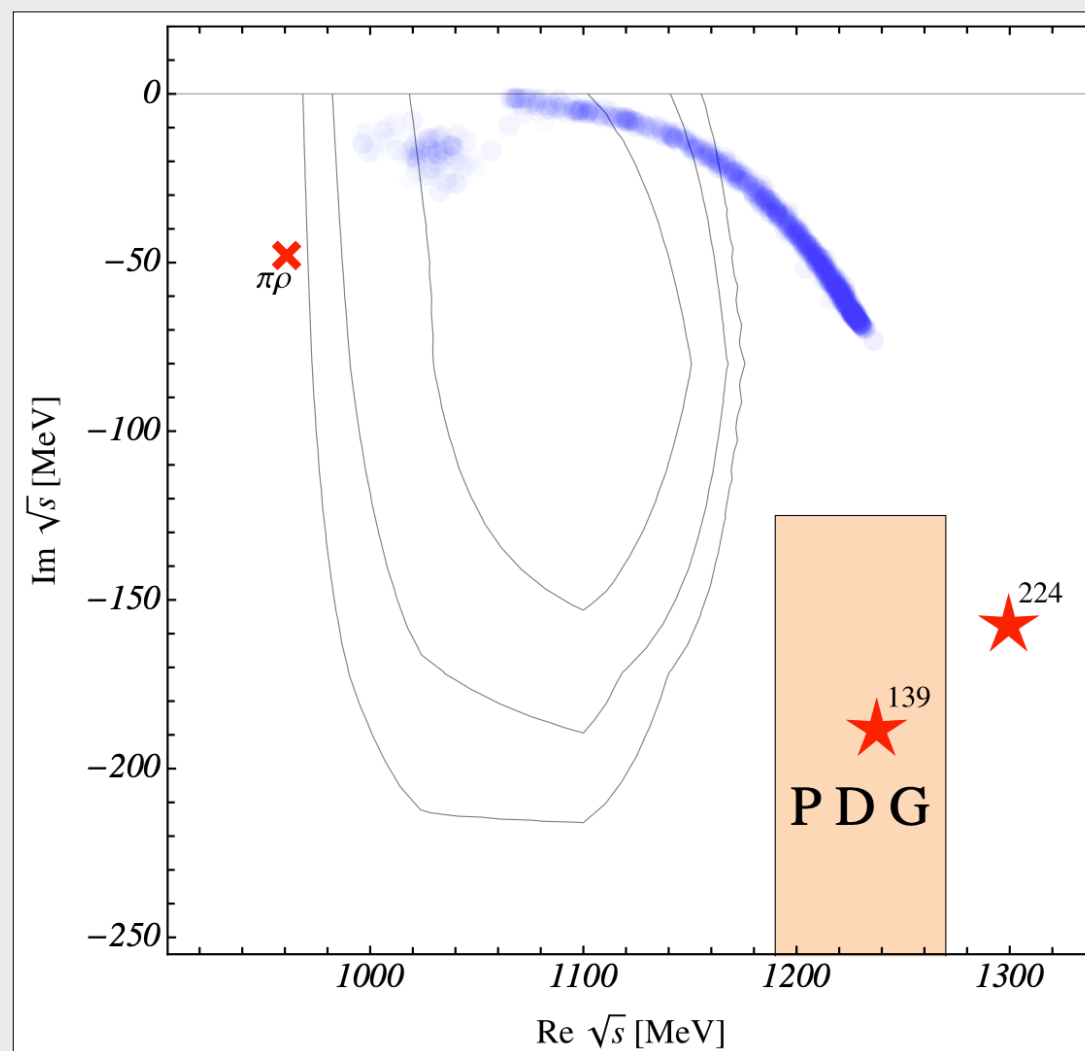


THE $a_1(1260)$ -RESONANCE FROM LATTICE QCD [2107.03973 \[hep-lat\]](https://arxiv.org/abs/2107.03973)



Maxim Mai, A. Alexandru, R. Brett, C. Culver
M. Döring, F. Lee, D. Sadasivan [GWQCD]



PHY-2012289



DE-SC0016582

DE-FG02-95ER40907

slides

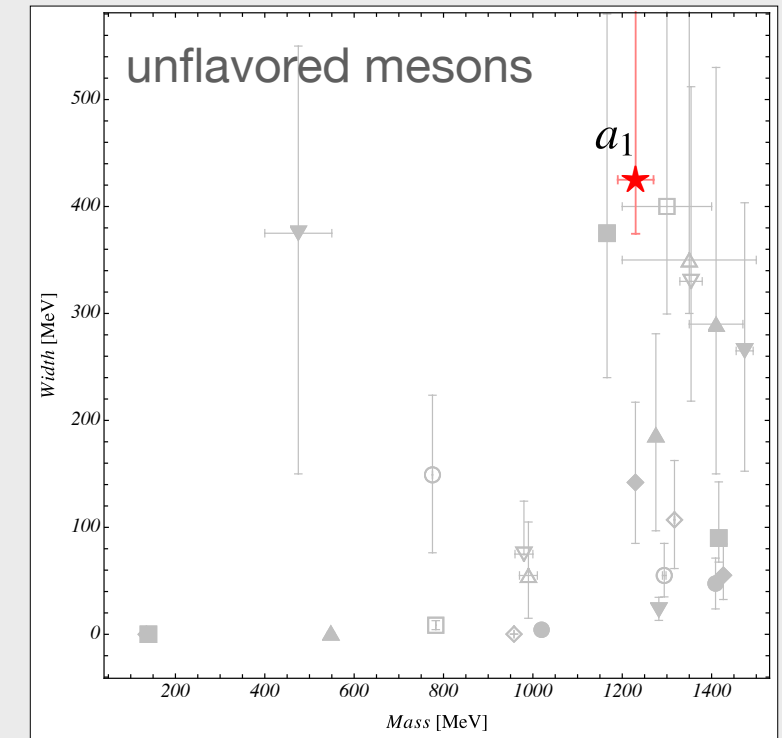
<https://maxim-mai.github.io/talks/LAT21-MM.pdf>

QCD SPECTRUM

Many states of QCD have large coupling to 3-body channels

- $\omega(782)$, $a_1(1260)$...
- exotic mesons: $\pi_1(1600)$, ... [exp. searches @ COMPASS, GlueX](#)
- Roper resonance $N^*(1440)$

This work: $a_1(1260)$ from lattice QCD



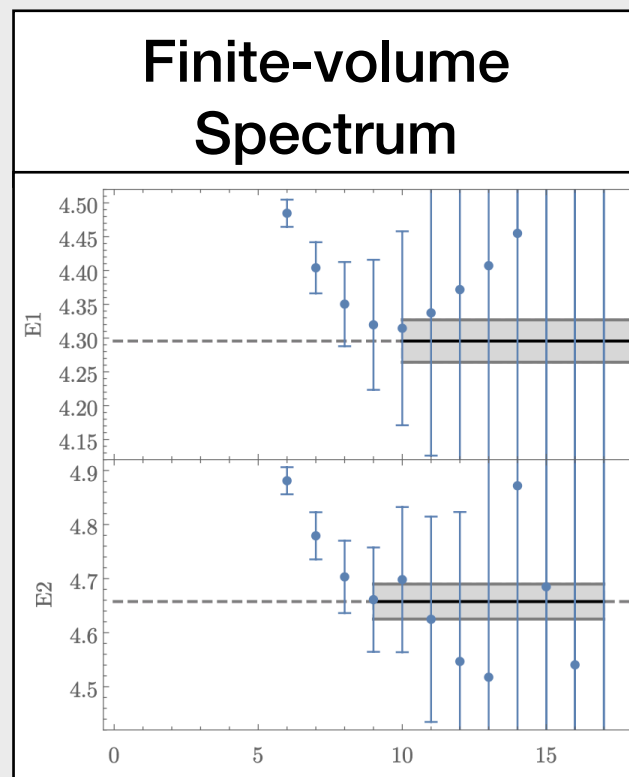
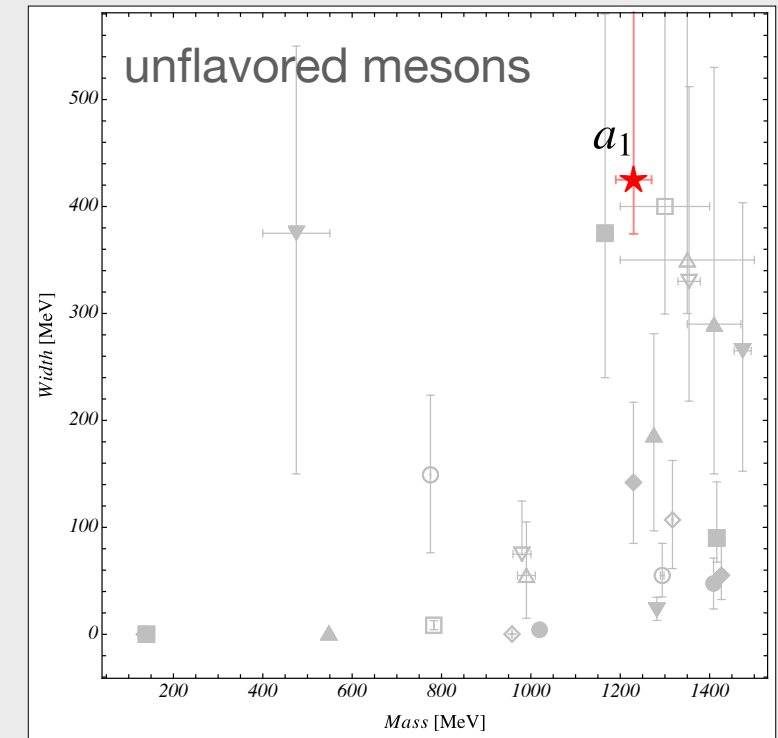
QCD SPECTRUM

Many states of QCD have large coupling to 3-body channels

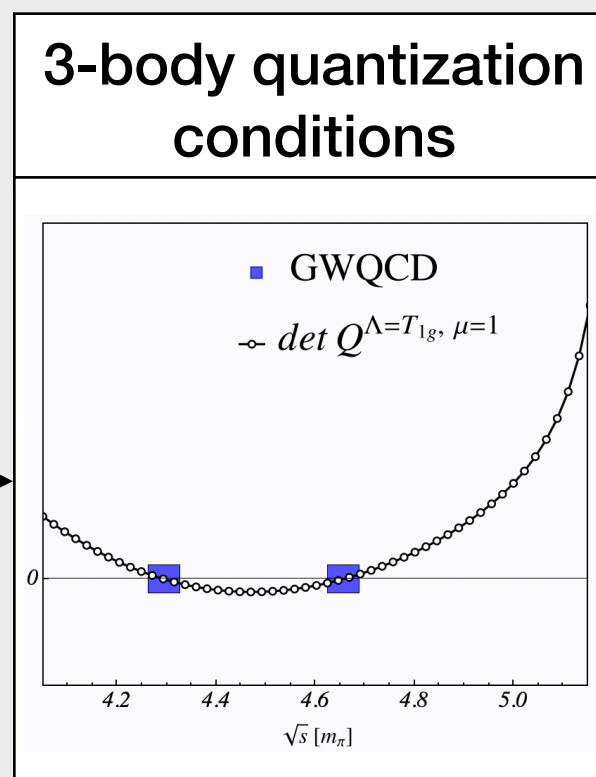
- $\omega(782)$, $a_1(1260)$...
- exotic mesons: $\pi_1(1600)$, ... [exp. searches @ COMPASS, GlueX](#)
- Roper resonance $N^*(1440)$

This work: **$a_1(1260)$ from lattice QCD**

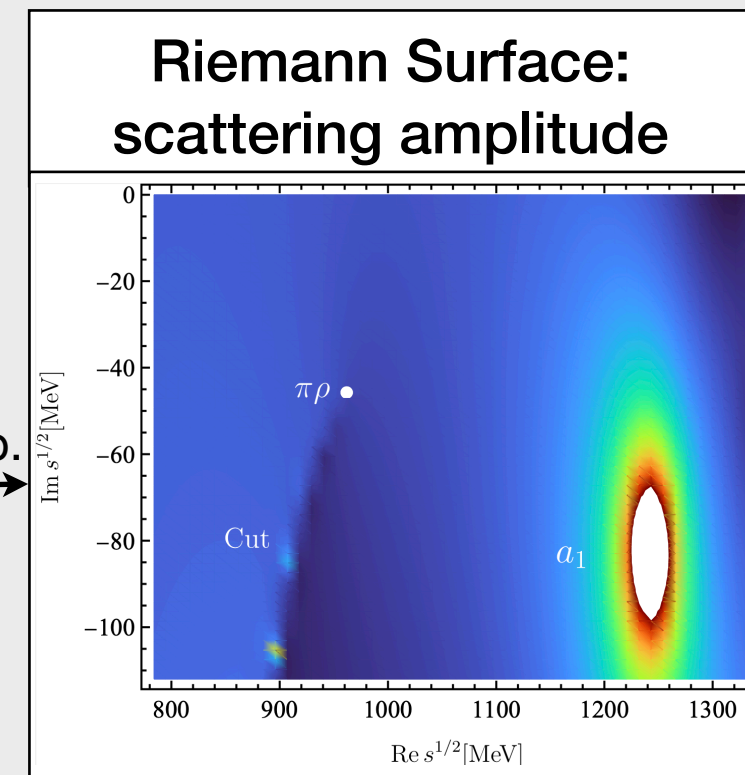
- Universal parameters from poles on the Riemann surface
- 3 step procedure:



energy eigenvalues



volume indep. quantities



FINITE-VOLUME SPECTRUM

GWQCD ensemble used for 2/3 pion calculations [Alexandru, Brett, Culver, Guo, Lee, Pelissier \(2013-2020\)](#)
[PRD87,PRD94,PRD98,PRD96,PRL117,PRD100](#)

Some key details: *(more in the next talk -- Ruairí Brett)*

- $N_f = 2$ dynamical fermions, LapH smearing
- $\mathbf{P}=(0,0,0)$, $m_\pi=224$ MeV, $m_\pi L=3.3$
- *GEVP with one-, two-, three-meson operators*
- *Relevant irrep(O_h) for $\mathbf{a}_1(1260)$ $I^G (J^P C) = 1^- (1^{++})$: T_{1g}*

Geometry	\mathbf{P}	Λ	$J^P (I^G = 1^-)$
Cubic	$\mathbf{P} = (0, 0, 0)$	T_{1g}	$1^+, 3^+, \dots$
		A_{1u}	$0^-, 4^-, \dots$

FINITE-VOLUME SPECTRUM

GWQCD ensemble used for 2/3 pion calculations [Alexandru, Brett, Culver, Guo, Lee, Pelissier \(2013-2020\)](#)
[PRD87,PRD94,PRD98,PRD96,PRL117,PRD100](#)

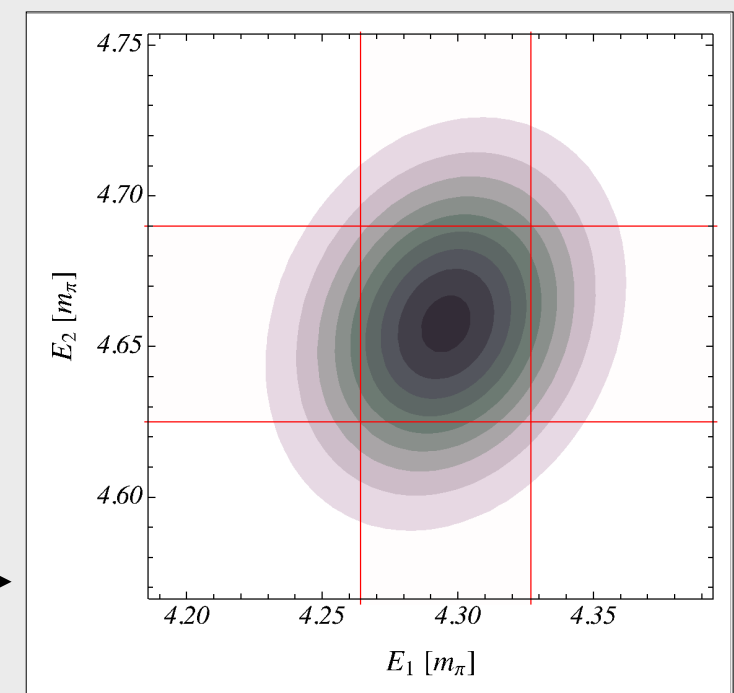
Some key details: *(more in the next talk -- Ruairí Brett)*

- $N_f = 2$ dynamical fermions, LapH smearing
- $\mathbf{P}=(0,0,0)$, $m_\pi=224$ MeV, $m_\pi L=3.3$
- *GEVP with one-, two-, three-meson operators*
- *Relevant irrep(O_h) for $a_1(1260)$ $I^G (J^P C) = 1^- (1^{++})$: T_{1g}*

Geometry	\mathbf{P}	Λ	$J^P (I^G = 1^-)$
Cubic	$\mathbf{P} = (0, 0, 0)$	T_{1g} A_{1u}	$1^+, 3^+, \dots$ $0^-, 4^-, \dots$

Key insights:

- 3-meson operators stabilize the excited state extraction
 c.f. need for $\rho\pi$ operators in pioneering 2-meson a_1 calculation [Lang et al. JHEP 04, 162 \(2014\)](#)
- high-momentum states are required: $\pi(0,0,0)\pi(1,1,0)\pi(-1,-1,0)$ etc..
- two interacting levels exists below 5π threshold \longrightarrow



3-BODY QUANTIZATION CONDITION

Discrete, real finite-volume (lattice) spectrum → continuous complex-valued amplitudes

- established in 2-body: Lüscher's method, extensions...
- 3-body methods matured (this session)

Lüscher, Gottlieb, Rummukainen, Feng, Li, Liu,
Döring, Briceño, Bernard, Meißner, Rusetsky...

Bedaque, Blanton, Briceño, Davoudi, Döring, Grieshammer, Guo,
Hammer, Hansen, MM, Meißner, Müller, Pang, Polejaeva, Romero-López,
Rusetsky, Sharpe, Wu

Reviews: Hansen/Sharpe(2019) MM/Döring/Rusetsky(2021)

3-BODY QUANTIZATION CONDITION

Discrete, real finite-volume (lattice) spectrum \rightarrow continuous complex-valued amplitudes

- established in 2-body: Lüscher's method, extensions...

Lüscher, Gottlieb, Rummukainen, Feng, Li, Liu, Döring, Briceño, Bernard, Meißner, Rusetsky...

- 3-body methods matured (this session)

Bedaque, Blanton, Briceño, Davoudi, Döring, Griesshammer, Guo, Hammer, Hansen, MM, Meißner, Müller, Pang, Polejaeva, Romero-López, Rusetsky, Sharpe, Wu

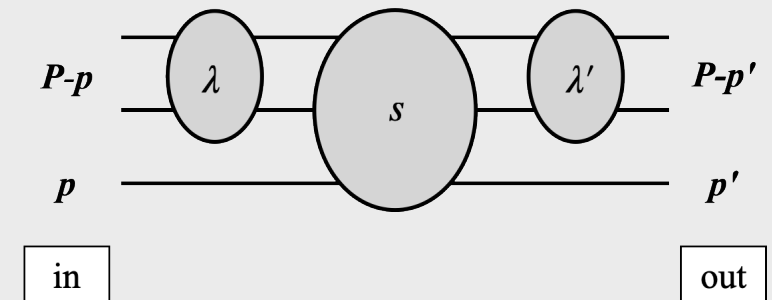
Reviews: Hansen/Sharpe(2019) MM/Döring/Rusetsky(2021)

Finite Volume Unitarity MM, Döring EPJA (2017) PRL (2019)

- basic idea:

(unitary three-body amplitude) $\int \frac{d^3\mathbf{k}}{(2\pi)^3} \rightarrow 1/L^3 \sum_k$ **(singular \Leftrightarrow three mesons are on-shell)** \Leftrightarrow **(energy eigenvalues)**

$$0 = \det \left[B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}^{\Lambda}$$



- extended to higher spin and coupled-channels: new degree of freedom (λ)

- ∞ -dim. determinant equation in $\mathbf{p} \in \frac{2\pi}{L} \mathbf{Z}^3 \rightarrow$ practical applications require truncation
 \rightarrow common to all quantization conditions

see discussion in e.g. MM/Döring/Rusetsky(2021)

"in the complexity it is only simplicity which can be interesting"

Steven Weinberg (May 3, 1933 - July 23, 2021)

"in the complexity it is only simplicity which can be interesting"

Steven Weinberg (May 3, 1933 - July 23, 2021)

3-BODY QUANTIZATION CONDITION (FVU)

$$0 = \det \left[B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}^{\Lambda}$$

"in the complexity it is only simplicity which can be interesting"

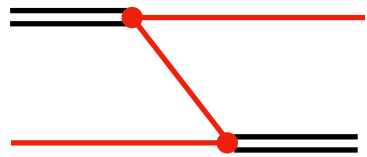
Steven Weinberg (May 3, 1933 - July 23, 2021)

3-BODY QUANTIZATION CONDITION (FVU)

$$0 = \det \left[B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}^{\Lambda}$$

one-particle exchange

- fixed by 3b-unitarity



- no free parameters

"in the complexity it is only simplicity which can be interesting"

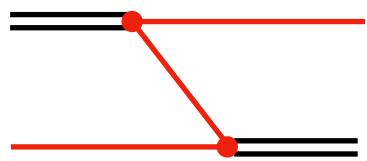
Steven Weinberg (May 3, 1933 - July 23, 2021)

3-BODY QUANTIZATION CONDITION (FVU)

$$0 = \det \left[B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}^{\Lambda}$$

one-particle exchange

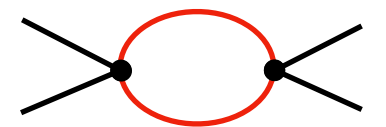
- fixed by 3b-unitarity



- no free parameters

two-body self-energy

- fixed by 2b-unitarity



- no free parameters

"in the complexity it is only simplicity which can be interesting"

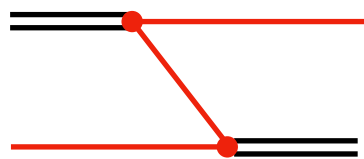
Steven Weinberg (May 3, 1933 - July 23, 2021)

3-BODY QUANTIZATION CONDITION (FVU)

$$0 = \det \left[B(s) + C(s) - 2L^3 E_p \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}^\Lambda$$

one-particle exchange

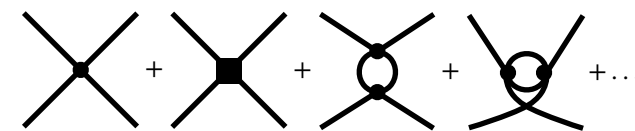
- fixed by 3b-unitarity



- no free parameters

two-body kernel

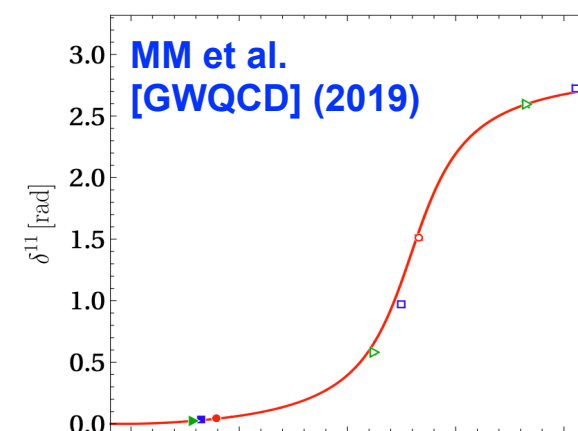
- dynamics of $l=1$ $\pi\pi$ system



- regular function \Rightarrow polynomial

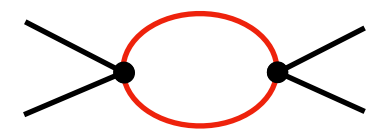
$$\tilde{K}_n^{-1}(s) = \sum_{i=0}^{n-1} a_i \cdot \sigma_p^i$$

- parameters (a_0, a_1) from cross-channel fit to $\pi\pi$ GWQCD levels



two-body self-energy

- fixed by 2b-unitarity



- no free parameters

"in the complexity it is only simplicity which can be interesting"

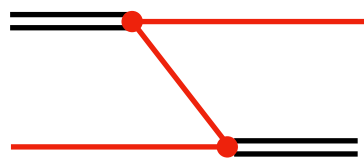
Steven Weinberg (May 3, 1933 - July 23, 2021)

3-BODY QUANTIZATION CONDITION (FVU)

$$0 = \det \left[B(s) + C(s) - 2L^3 E_p \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}^\Lambda$$

one-particle exchange

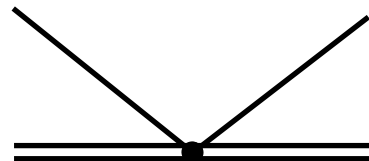
- fixed by 3b-unitarity



- no free parameters

three-body force

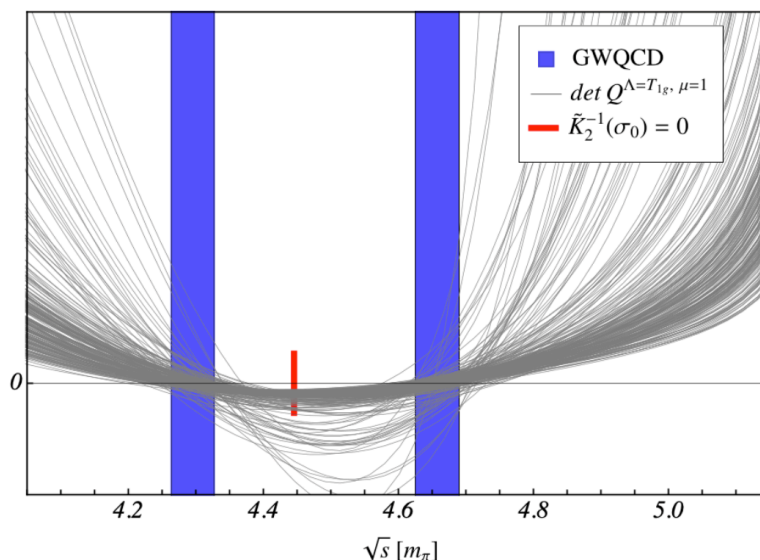
- dynamics of $\rho\pi$ system



- regular function \Rightarrow Laurent series

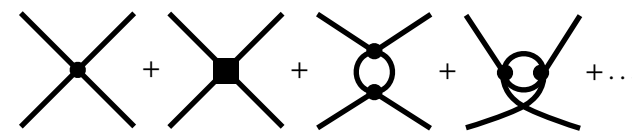
$$C_{\ell'\ell}(s, \mathbf{p}', \mathbf{p}) = \sum_{i=-1}^{\infty} c_{\ell'\ell}^{(i)}(\mathbf{p}', \mathbf{p})(s - \mathbf{m}_{a_1}^2)^i$$

- fit to 3-body levels



two-body kernel

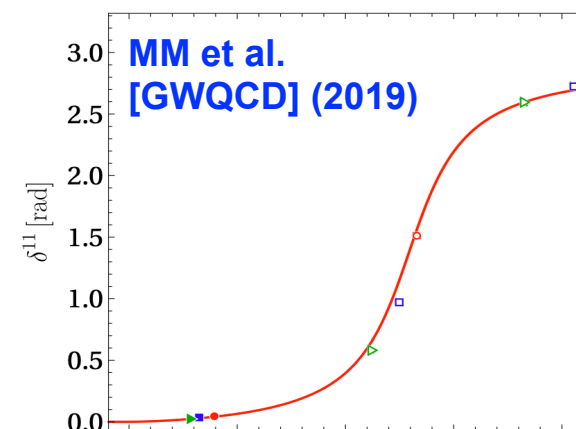
- dynamics of $l=1$ $\pi\pi$ system



- regular function \Rightarrow polynomial

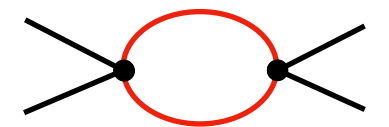
$$\tilde{K}_n^{-1}(s) = \sum_{i=0}^{n-1} a_i \cdot \sigma_p^i$$

- parameters (a_0, a_1) from cross-channel fit to $\pi\pi$ GWQCD levels



two-body self-energy

- fixed by 2b-unitarity



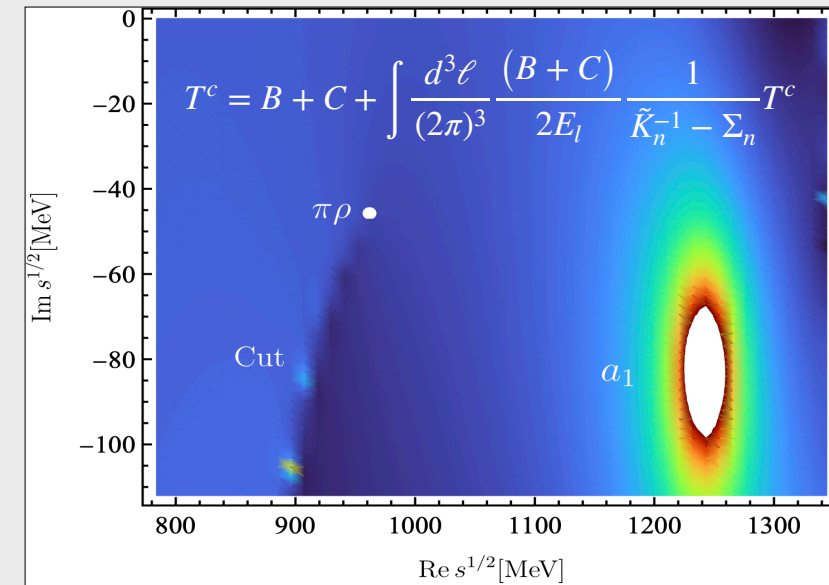
- no free parameters

RESULTS

Resonance poles

- ∞ -vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term

- best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{s - m_{a_1}^2} g_{\ell} |\mathbf{p}|^{\ell} + c \delta_{\ell'0} \delta_{\ell 0}$
 ...with large correlations

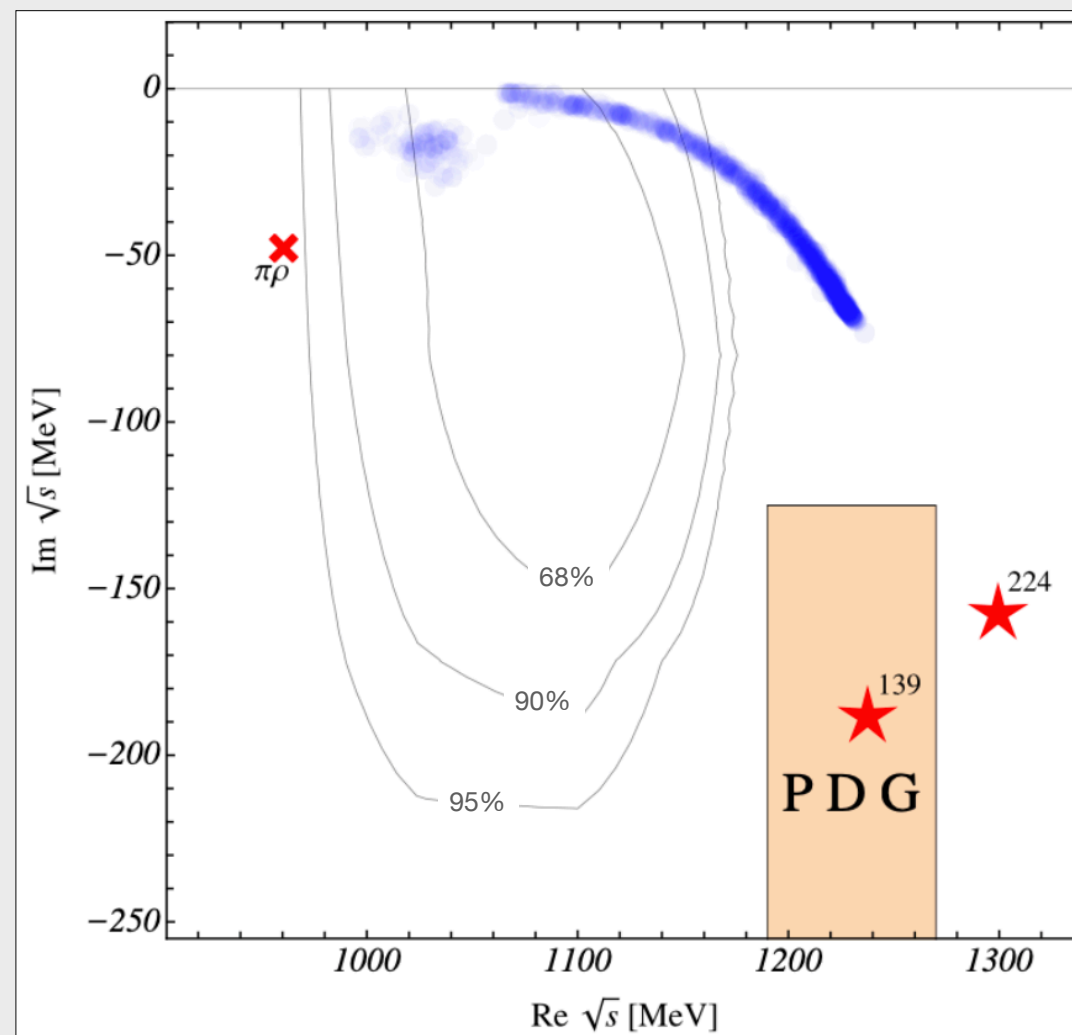
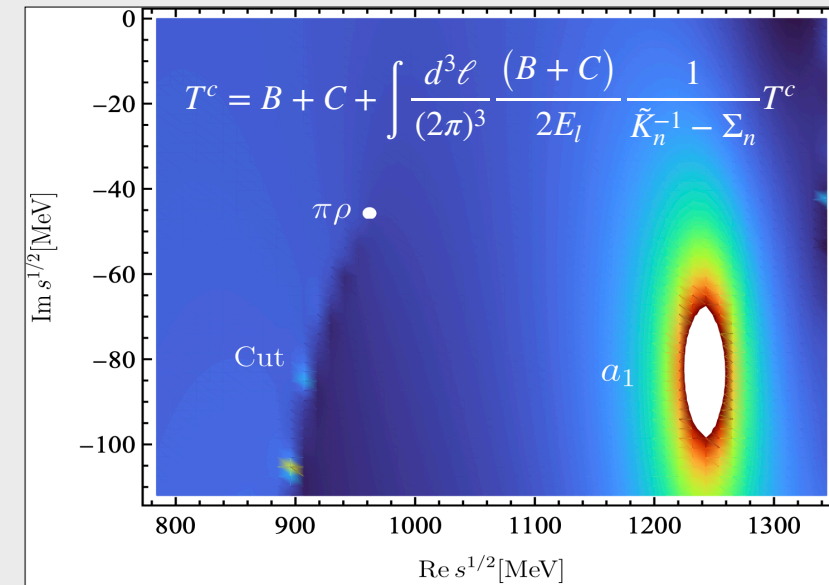


RESULTS

Resonance poles

- ∞ -vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term

- best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{s - m_{a_1}^2} g_{\ell} |\mathbf{p}|^{\ell} + c \delta_{\ell'0} \delta_{\ell 0}$
 ...with large correlations

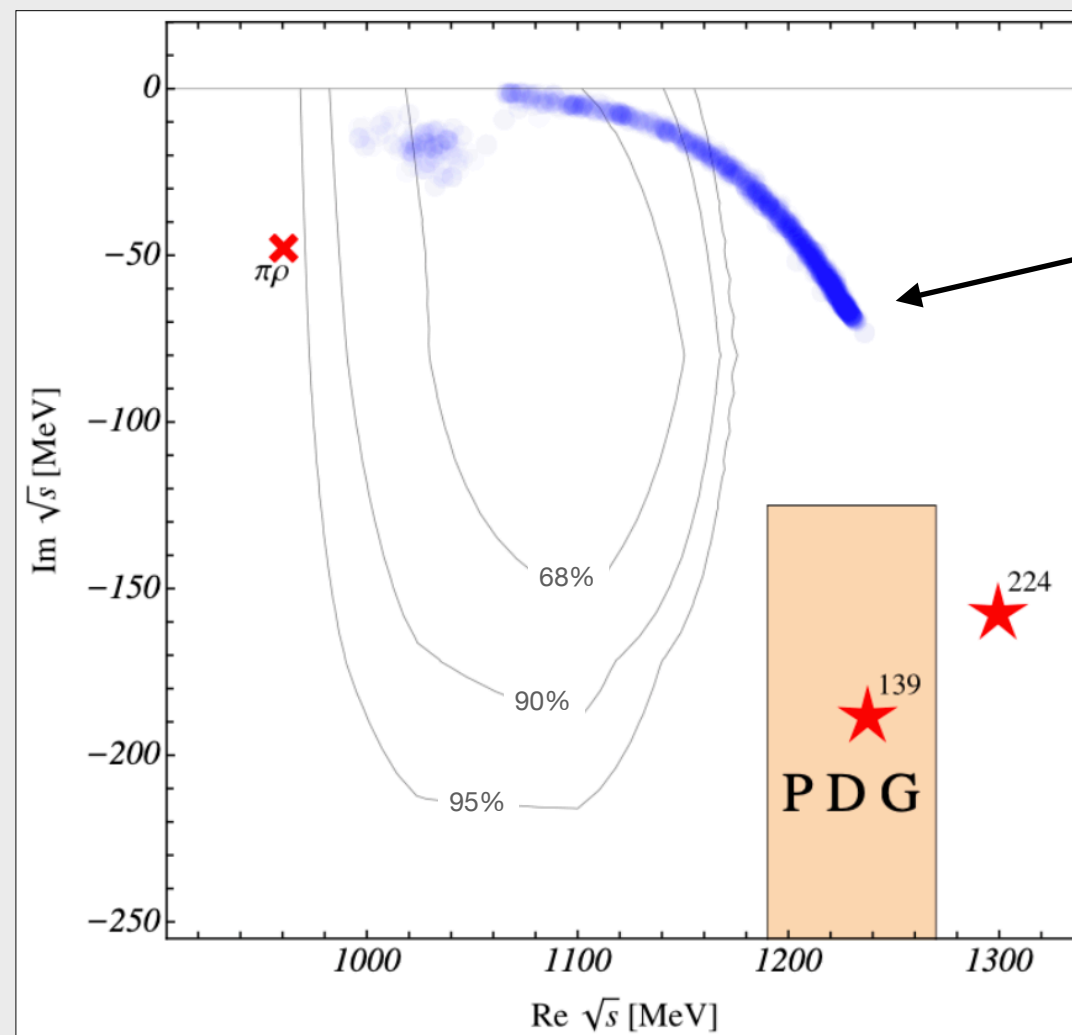
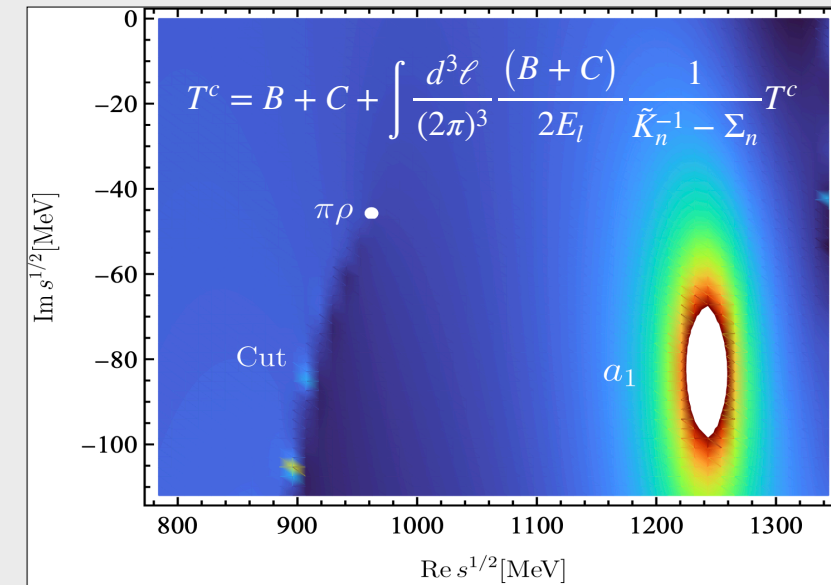


RESULTS

Resonance poles

- ∞ -vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term

- best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{s - m_{a_1}^2} g_{\ell} |\mathbf{p}|^{\ell} + c \delta_{\ell'0} \delta_{\ell 0}$
 ...with large correlations



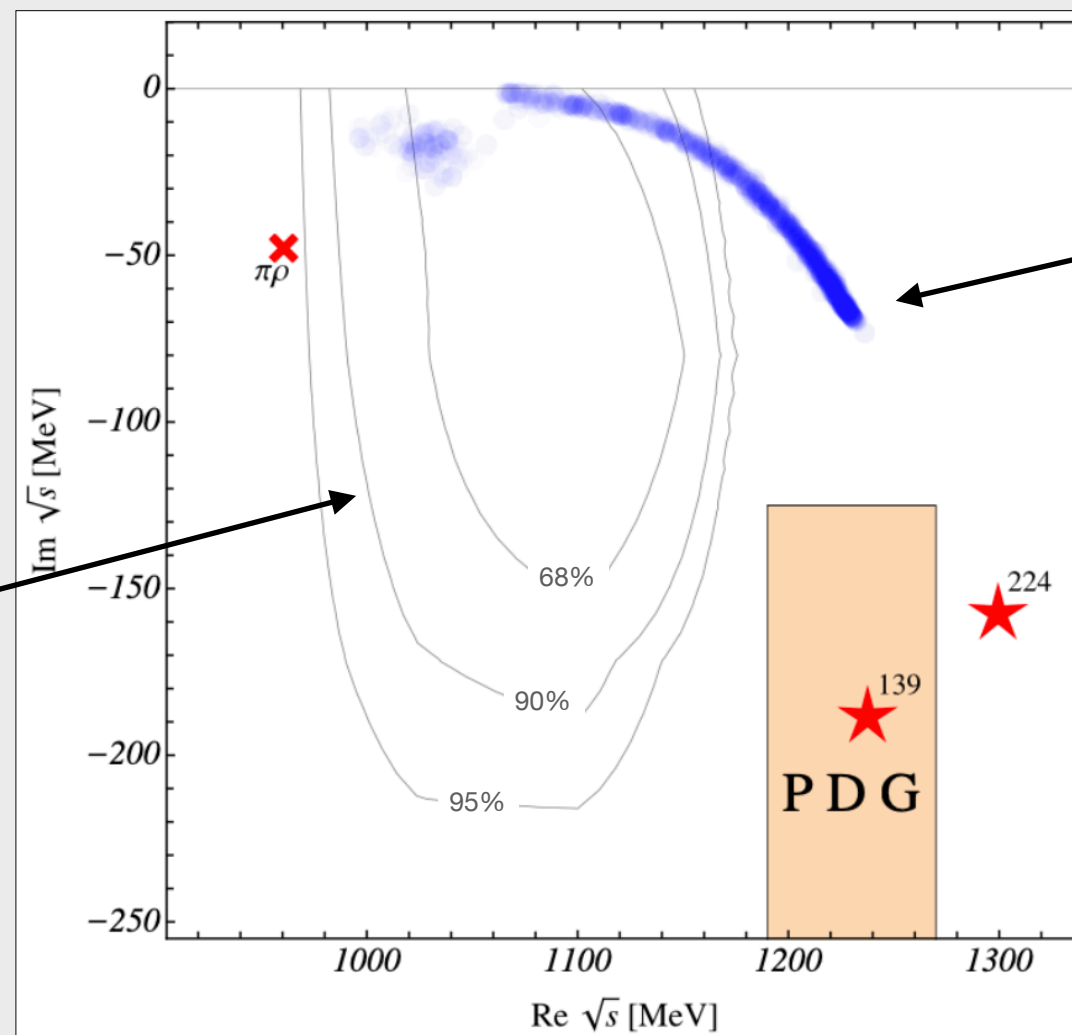
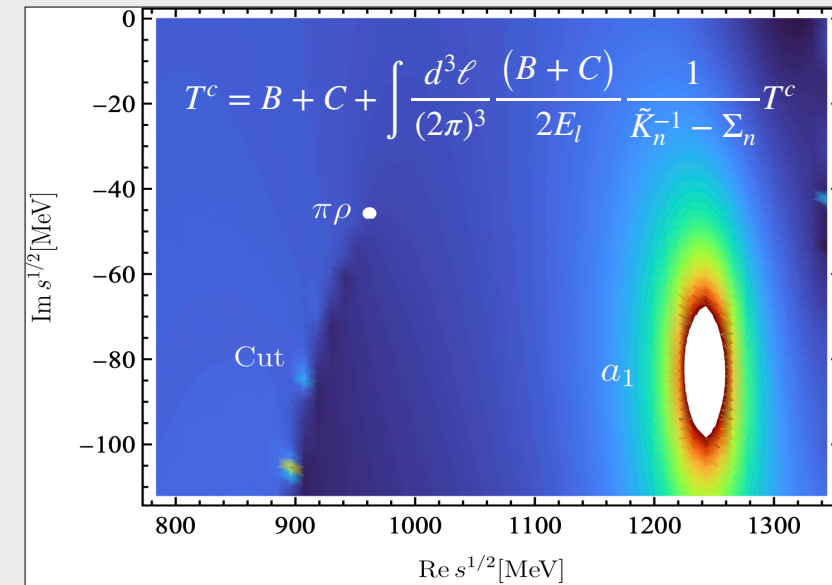
distribution is finite
 $\sim 1230 - i60$ MeV

RESULTS

Resonance poles

- ∞ -vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term

- best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{s - m_{a_1}^2} g_{\ell} |\mathbf{p}|^{\ell} + c \delta_{\ell'0} \delta_{\ell 0}$
 ...with large correlations



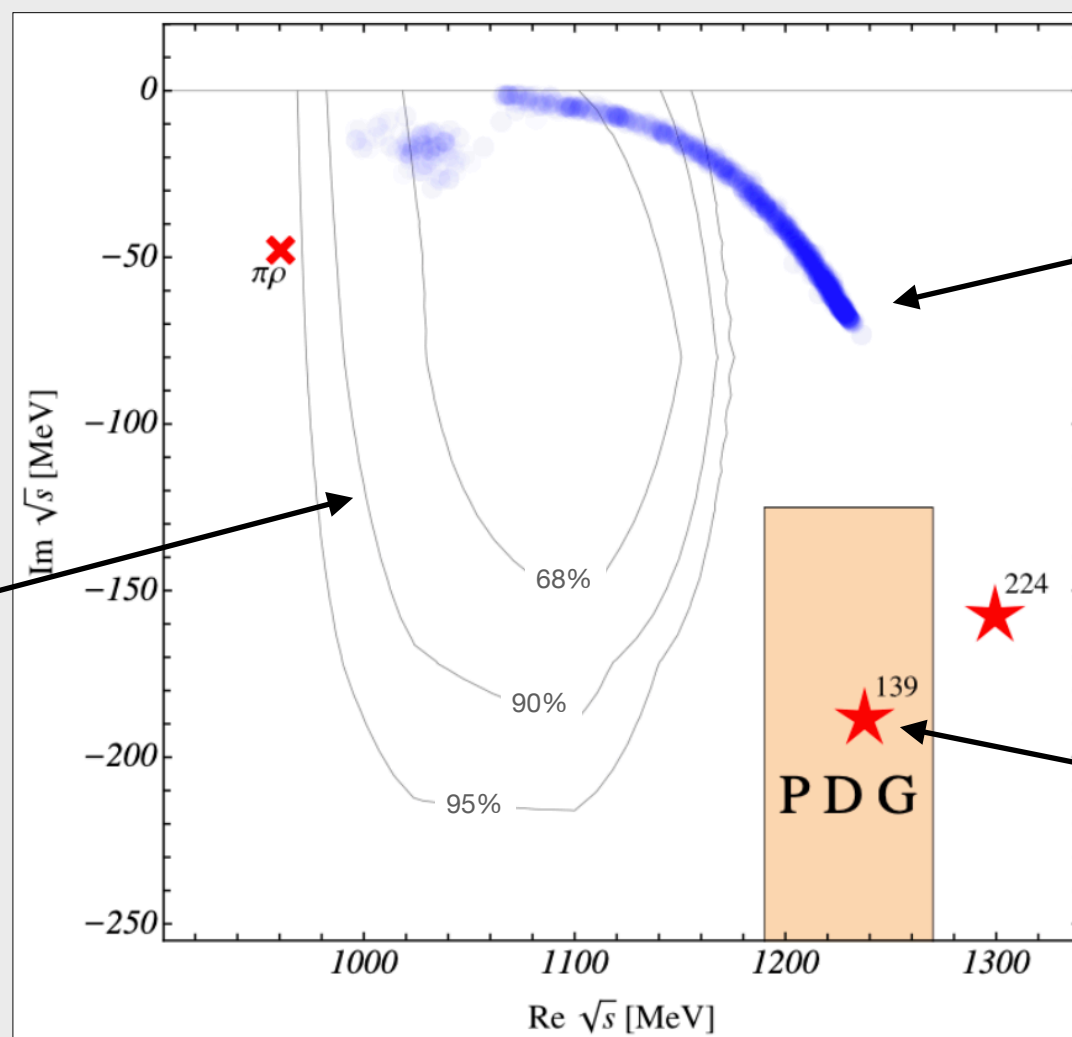
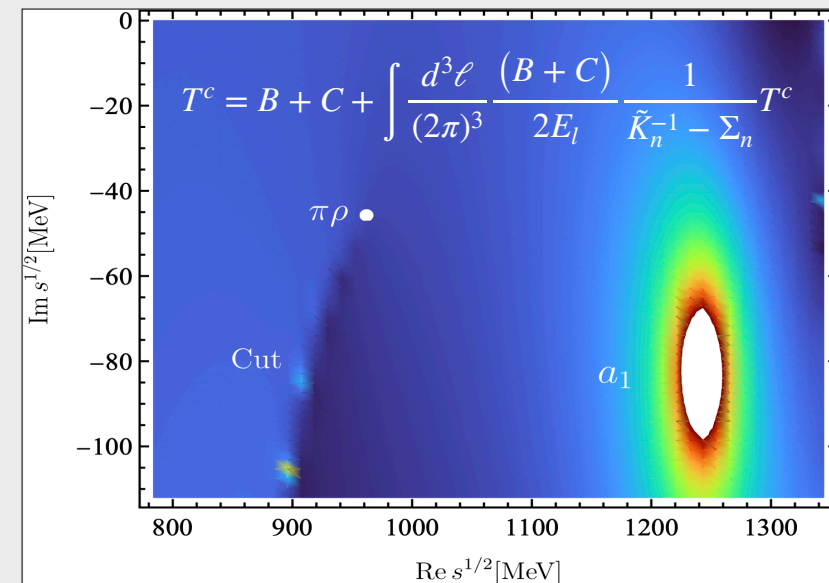
distribution is finite
 $\sim 1230 - i60$ MeV

Lüscher + Breit-Wigner approximation
 c.f. Lang et al. JHEP 04(2014)

RESULTS

Resonance poles

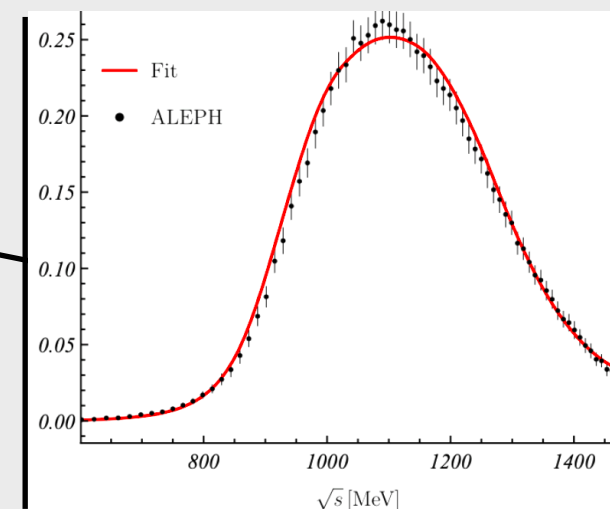
- ∞ -vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{s - m_{a_1}^2} g_{\ell} |\mathbf{p}|^{\ell} + c \delta_{\ell'0} \delta_{\ell 0}$
...with large correlations



distribution is finite
~1230-i60 MeV

Lüscher + Breit-Wigner approximation

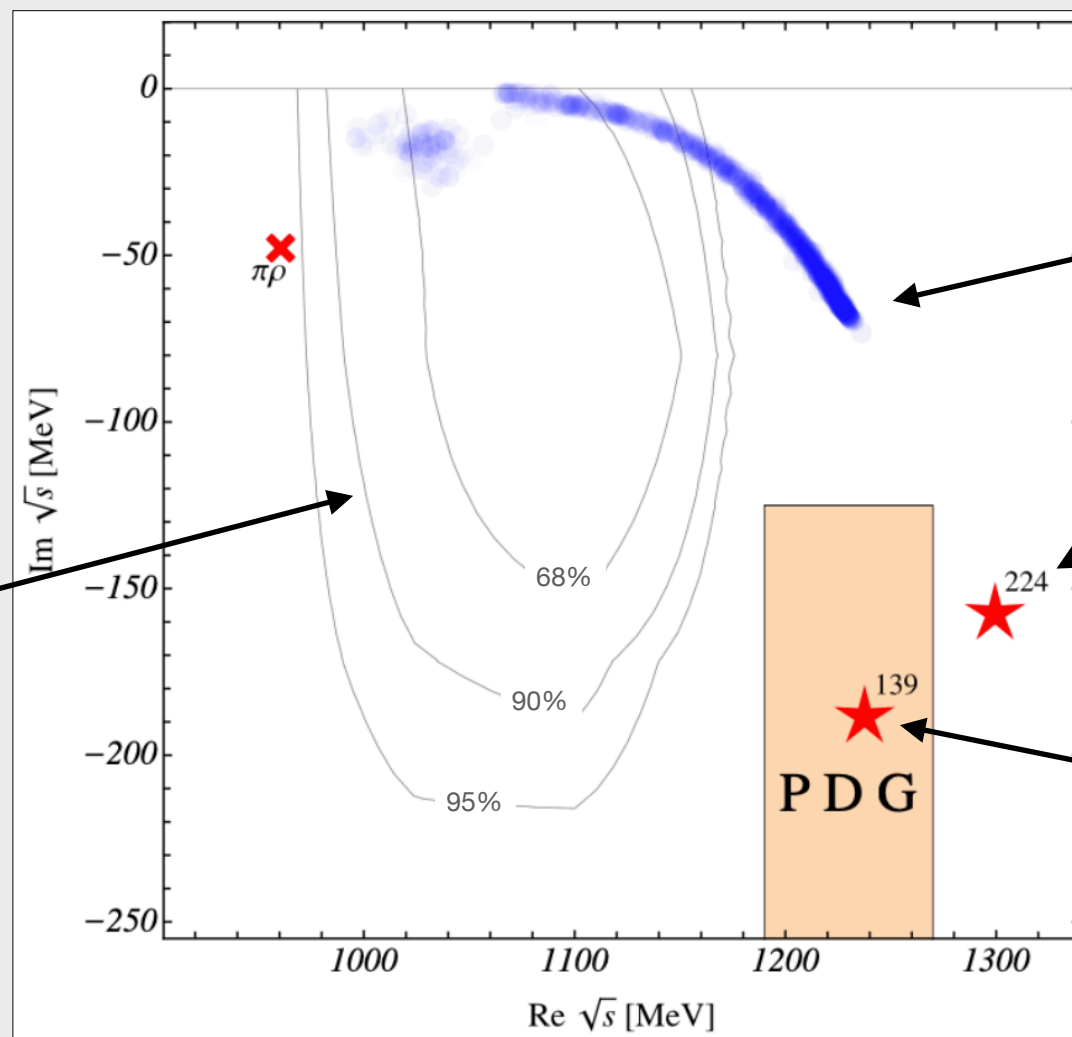
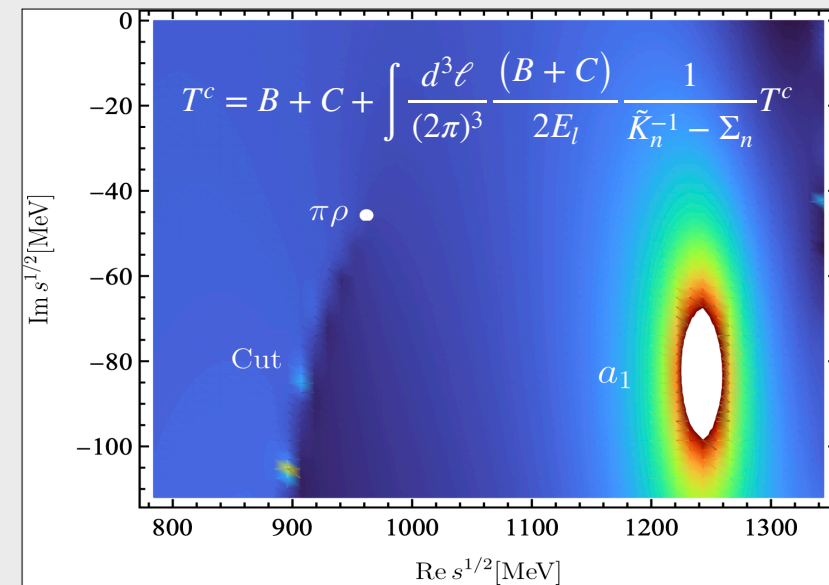
c.f. Lang et al. JHEP 04(2014)



RESULTS

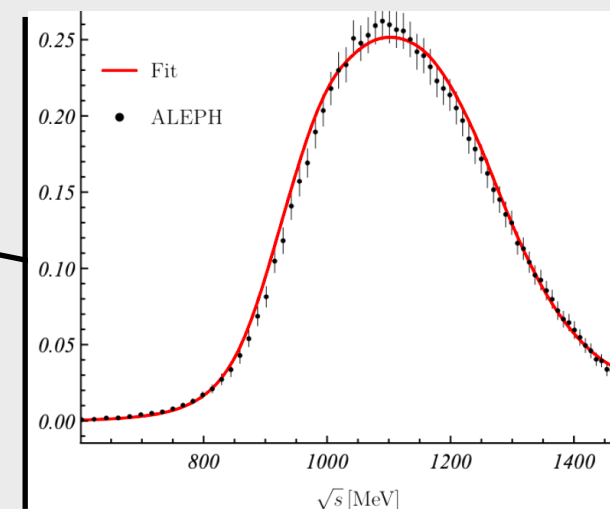
Resonance poles

- ∞ -vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{s - m_{a_1}^2} g_{\ell} |\mathbf{p}|^{\ell} + c \delta_{\ell'0} \delta_{\ell 0}$
...with large correlations



distribution is finite
~1230-i60 MeV

(naive) chiral extrapolation
confirms expectation



Lüscher + Breit-Wigner
approximation

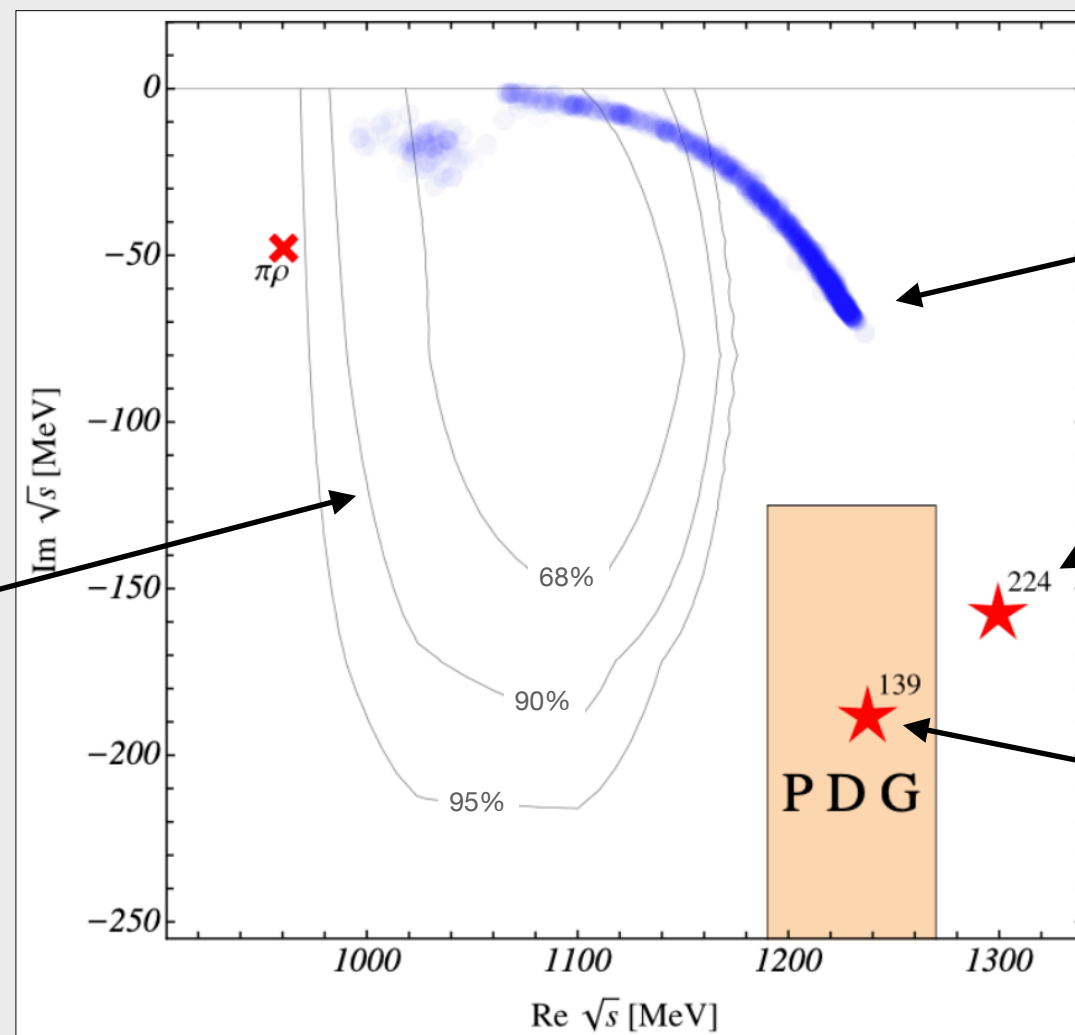
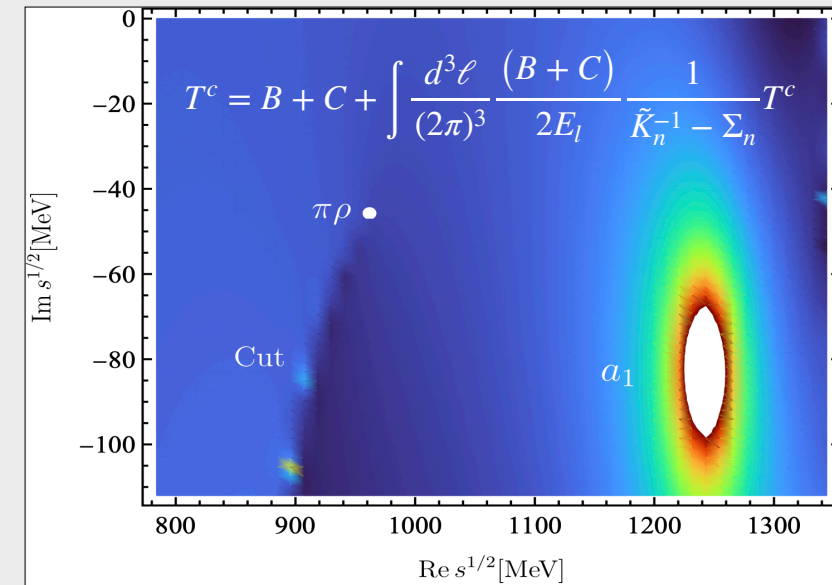
c.f. Lang et al. JHEP 04(2014)

RESULTS

Resonance poles

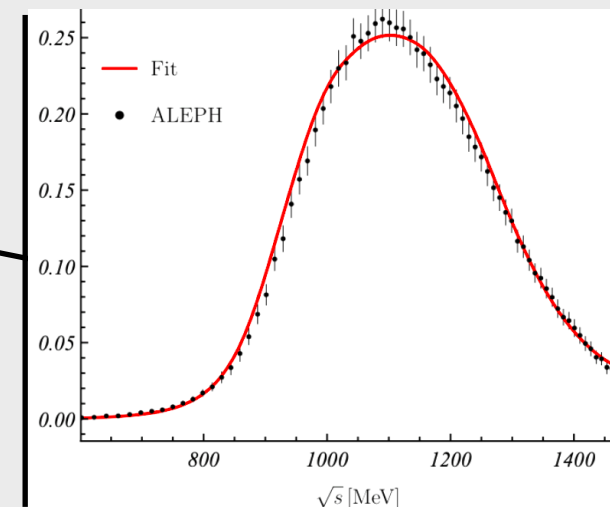
- ∞ -vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term

- best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{s - m_{a_1}^2} g_{\ell} |\mathbf{p}|^{\ell} + c \delta_{\ell'0} \delta_{\ell 0}$
 ...with large correlations



distribution is finite
 $\sim 1230 - i60$ MeV

(naive) chiral extrapolation
 confirms expectation



Lüscher + Breit-Wigner
 approximation

c.f. Lang et al. JHEP 04(2014)

THANK YOU